HELGASON-MARCHAUD INVERSION FORMULAS FOR RADON TRANSFORMS ON CONSTANT CURVATURE SPACES

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Abstract. Let \( \tilde{f}(\xi) \) be the \( k \)-dimensional totally geodesic Radon transform of a function \( f \) on the real hyperbolic space \( \mathbb{H}^n \), \( 1 \leq k \leq n - 1 \). By averaging \( \tilde{f}(\xi) \) over all \( \xi \) at a distance \( \theta \) from \( x \in \mathbb{H}^n \), and applying Riemann-Liouville fractional differentiation in \( \theta \), Helgason has recovered \( f(x) \). This method blows up if \( f \) does not decrease sufficiently fast. The situation can be saved if one employs Marchaud’s fractional derivatives instead of the Riemann-Liouville ones. New inversion formulas for \( \tilde{f}(\xi) \), \( f \in L^p \), are obtained. Similar problems are studied for geodesic transforms on the sphere \( S^n \), where the theory is qualitatively different.

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1. Introduction

Following [6, p. 96], let $X$ be either the hyperbolic space $\mathbb{H}^n$ or the sphere $S^n$; $\Xi$ is the set of all $k$-dimensional totally geodesic submanifolds $\xi \subset X$, $1 \leq k \leq n-1$. For $\theta > 0$ we denote

$$
\dot{f}_\theta(\xi) = \int_{d(x,\xi)=\theta} f(x)dm(x), \quad \varphi(\xi) = \int_{d(x,\xi)=\theta} \varphi(x)d\mu(\xi),
$$

where $d(\cdot, \cdot)$ designates a geodesic distance, $d\mu(x)$ and $dm(\xi)$ are the relevant normalized measures. For $\theta = 0$ operators (1.1) coincide with the Radon transform $\dot{f}(\xi)$ and its dual $\varphi(\xi)$. Helgason [5, 6, p. 97] has proved the following formulae:

\begin{align*}
(1.2) \quad f(x) &= c \left[ \left( \frac{d}{d(u^2)} \right)^k \int_0^u (\dot{f})^\nu_{\cosh^{-1}(\nu^{-1})}(x)(u^2 - v^2)^{k/2-1}dv \right]_{u=1}, \quad X = \mathbb{H}^n; \\
(1.3) \quad f(x) &= \frac{c}{2} \left[ \left( \frac{d}{d(u^2)} \right)^k \int_0^u (\dot{f})^\nu_{\cos^{-1}(\nu)}(x)v^k(u^2 - v^2)^{k/2-1}dv \right]_{u=1}, \quad X = S^n.
\end{align*}

Here $f \in C^\infty_c(\mathbb{H}^n)$ or $f \in C^\infty(S^n)$, $c^{-1} = (k-1)!/\sigma_k/2^{k+1}$, $\sigma_k$ is the area of the unit sphere $S^k \subset \mathbb{R}^{k+1}$.

**QUESTION:** Are (1.2) and (1.3) applicable to $f \in L^p(X)$?

For (1.3) the answer is “yes” [9] provided that all derivatives are understood in the $L^p$-norm or in the a.e. sense. In the noncompact case (1.2) the answer is in general “no”, as we shall see below. It is known [3, 9] that for $f \in L^p(X)$, $\dot{f}(\xi)$ exists for almost all $\xi \in \Xi$ provided that $1 \leq p \leq \infty$ if $X = S^n$, and

$$
1 \leq p < (n-1)/(k-1)
$$

if $X = \mathbb{H}^n$. The last condition is sharp. The point is that the integral in (1.2) can be divergent even if $\dot{f}$ is well defined. Discrepancies of this kind are usual in fractional calculus [8] where implementation of Marchaud’s method saves situation. The following statements demonstrate how this method works in our case for $k = 1$.

**Theorem A.** Let $f \in L^p(\mathbb{H}^n)$, $1 \leq p < \infty$. Then

$$
f = \frac{1}{\pi} \int_0^\infty \frac{(\dot{f})^\nu - (\dot{f})^\nu_0}{\sinh^2 \theta} \cosh \theta d\theta, \quad \int_0^\infty = \lim_{\varepsilon \to 0} \int_\varepsilon^\infty,
$$

where the limit is understood in the $L^p$-norm and in the a.e. sense. If $f \in C_0(\mathbb{H}^n) \cap L^p(\mathbb{H}^n)$, where $C_0(\mathbb{H}^n)$ is the space of continuous functions vanishing at infinity, then $\lim_{\varepsilon \to 0}$ can be treated in the sup-norm.
Theorem B. Let $f \in L^p(S^n)$, $1 \leq p < \infty$. Then

\begin{equation}
 f = \frac{(\hat{f})^\vee}{2\pi} + \frac{1}{2\pi} \int_0^{\pi/2} \frac{(\hat{f})^\vee - (\hat{f})_0^\vee}{\sin^2 \theta} \cos \theta \, d\theta, \quad \int_0^{\pi/2} = \lim_{\varepsilon \to 0} \int_{\varepsilon}^{\pi/2},
\end{equation}

where the limit is understood in the $L^p$-norm and in the a.e. sense. If $f \in C(S^n)$, then $\lim_{\varepsilon \to 0}$ is uniform.

Analogous formulae for $k > 1$ involve finite differences of higher orders or/and derivatives which must be understood in a certain $L^q$-norm or in the a.e. sense (see Theorems 2.2, 2.4, and 3.3 below, and also [9; Section 4.2]).

Historical remark. Apart from (1.2) and (1.3), Helgason [4] suggested another inversion formula $f = P_{k/2}(\Delta)(\hat{f})^\vee$, in which $k$ is even, $f$ is smooth, and $P_{k/2}(\Delta)$ is a polynomial of the Laplace-Beltrami operator $\Delta$. This formula was extended to all $k$ by Berenstein and Casadio Tarabusi (for $X = \mathbb{H}^n$) in [1] and by the author (for $X = S^n, \mathbb{H}^n$) in [9, 10], using different approaches. Investigation of Radon transforms on $L^p$-spaces was initiated by Oberlin and Stein [7], and Strichartz [13]. For $f \in L^p$, $\hat{f}$ can be inverted explicitly by making use of continuous wavelet transforms and the convolution-backprojection method [2, 3, 9, 11]. Applicability of (1.2) to $f \in L^p$ was not studied before. We complete this gap below. References to the relevant important works by S. Gindikin, E. Grinberg, S. Ishikawa, T. Kakehi, A. Kurusa, E.T. Quinto and others can be found in [2, 3, 6, 9-11].

2. Inversion formulae of the Marchaud type

2.1. The case $X = \mathbb{H}^n$. We start with the following equality

\begin{equation}
 \sigma_{k-1} \int_0^{\infty} F(\cosh \theta \cosh r) \sinh^{k-1} r dr = \hat{F}(\cosh \theta),
\end{equation}

which is due to Helgason [5, 6]. Here

\begin{equation}
 F(t) = (M_t f)(x) = \frac{\sigma_{n-1}}{(t^2 - 1)^{(1-n)/2}} \int_{y \in \mathbb{H}^n : [x, y] = t} f(y) d\sigma(y), \quad [x, y] = -x_1 y_1 - \ldots - x_n y_n + x_{n+1} y_{n+1}; \quad \hat{F}(\cosh \theta) = (\hat{f})_0^\vee(x).
\end{equation}

From this point our argument becomes different. The first question is whether (2.1) is true for $f \in L^p$. By setting $u = \cosh \theta$, $v = \cosh \theta \cosh r$, we write (2.1) as

\begin{equation}
 \sigma_{k-1} \int_u^{\infty} F(v)(v^2 - u^2)^{k/2 - 1} dv = u^{k-1} \hat{F}(u).
\end{equation}
The left hand side of (2.3) represents a hyperbolic convolution of the form

\[(k \ast f)(x) = \int_X f(y)k(x,y)dy\]

with \(k(\tau) = \sigma_{k-1}\sigma_{n-1}(\tau^2 - u^2)^{k/2-1}(\tau^2 - 1)^{1-n/2}\) (we use a standard notation \(a^\lambda = a^\lambda\) if \(a > 0\) and 0 otherwise). By Young’s inequality

\[||k \ast f||_q \leq ||f||_p||k||_r, \quad 1 \leq p \leq q \leq \infty, \quad 1 - \frac{1}{p} + \frac{1}{q} = \frac{1}{r},\]

the left hand side of (2.3) (written as (2.4)) is a linear bounded operator from \(L^p\) to \(L^q\) provided \(0 \leq q^{-1} \leq p^{-1} - (k - 1)/(n - 1)\). This condition agrees with (1.4). Thus for all \(f \in L^p\) with \(p\) in the maximal range \(1 \leq p < (n - 1)/(k - 1)\), the left hand side of (2.3) is finite for all \(u > 0\) and almost all \(x\). This justifies all operations having been used in [5, 6] for derivation of (2.1).

In order to recover \(f\) we write (2.3) in the form

\[(I^{k/2}_- g_x)(u) = \psi_x(u),\]

\[(g_x(v) = v^{-1/2}(M_\sigma f)(x), \quad \psi_x(u) = \frac{u^{(k-1)/2}}{\pi^{k/2}} \frac{1}{\sqrt{v^2 - u}} \cos^{-1}\sqrt{\pi}(x),\]

where \((I^\alpha_\sigma g)(u) = \frac{1}{\Gamma(\alpha)} \int_0^\infty g(v)(v-u)^{\alpha-1}dv\) denotes the Weyl fractional integral of order \(\alpha\) [8, 12]. A simple calculation yields

\[(I^\alpha_\sigma g_x)(u) = \frac{2}{\Gamma(\alpha)\sigma_{n-1}} \int_X f(y)(|x|, |y|^2 - u)^{\alpha-1}(|x|, |y|^2 - 1)^{1-n/2}dy,\]

and by (2.5), for \(1 \leq p \leq \infty,\)

\[||I^\alpha_\sigma g_x||_q \leq c||f||_p, \quad 0 \leq \frac{1}{q_\alpha} < 1 - \frac{2\alpha - 1}{n - 1}\]

The condition

\[(2\alpha - 1)/(n - 1) < 1/p \leq 1\]

is necessary for the existence of (2.8). It means that implementation of the semigroup property \(I^\alpha_\sigma I^\beta_\sigma = I^{\alpha+\beta}_\sigma\) (as in [5, 6]) may lead to divergent integrals.
Example 2.1. For \( k = 1 \) the Radon transform \( \hat{f} \), \( f \in L^p \), is well defined for all \( p \in [1, \infty) \). Formally \( I_{-}^{-1/2}I_{-}^{1/2}g_x = I_{-}^{1}g_x \). By (2.10) with \( \alpha = 1 \) the integral \( I_{-}^{1}g_x \) can be divergent if \( p \geq n - 1 \) (the case \( n = 2 \) is especially unpleasant).

In order to avoid unnatural restrictions on \( k, p \) and \( n \), we invert (2.6) with the aid of Marchaud’s fractional derivative (see [8, Section 10] for the general background of Marchaud’s method). Let \( \alpha = k/2 \). Given a positive integer \( \ell > \alpha \), we denote (cf. [8, p. 157])

\[
(\mathbb{D}_x^\alpha \psi)(u) = \frac{1}{\varphi_\ell(\alpha)} \int_0^\infty \left[ \sum_{j=0}^{\ell} \binom{\ell}{j} (-1)^j \psi(u + j t) \right] \frac{dt}{t^{1+\alpha}}, \quad \varepsilon > 0,
\]

\[
\varphi_\ell(\alpha) = \int_0^\infty (1 - e^{-t})\ell t^{1-\alpha} dt
\]

\[
= \begin{cases} 
\Gamma(-\alpha) \sum_{j=1}^{\ell} \binom{\ell}{j} (-1)^j j^\alpha, & \alpha \neq 1, 2, \ldots, \ell - 1, \\
\frac{(-1)^{1+\alpha}}{\alpha!} \sum_{j=1}^{\ell} \binom{\ell}{j} (-1)^j j^\alpha \log j, & \alpha = 1, 2, \ldots, \ell - 1.
\end{cases}
\]

The (right-sided) Marchaud’s fractional derivative of \( \psi \) of order \( \alpha \) is defined by

\[
\mathbb{D}_x^\alpha \psi = \lim_{\varepsilon \to 0} \mathbb{D}_x^\alpha \varepsilon \psi [8, 12].
\]

Theorem 2.2. Let \( \varphi = \hat{f} \), \( f \in L^p(\mathbb{H}^n) \), \( (k - 1)/(n - 1) < 1/p \leq 1 \). Then

\[
f(x) = \frac{\pi^{-k/2}}{\varphi_\ell(k/2)} \int_0^\infty \left[ \sum_{j=0}^{\ell} \binom{\ell}{j} (-1)^j (1 + j t)^{(k-1)/2} \varphi^{-1}_{\varepsilon} \left( \frac{\sqrt{1 + j t}}{\varepsilon} \right) \right] \frac{dt}{t^{1+k/2}},
\]

\[
\int_0^\infty = \lim_{\varepsilon \to 0} \int_0^\infty,
\]

where the limit is understood in the \( L^p \)-norm and in the a.e. sense. If \( f \in C_0 \cap L^p \),
then \( \lim_{\varepsilon \to 0} \) can be treated in the sup-norm.

Remark 2.3. The formula (2.12) is a “regularized” version of Helgason’s formula (1.2) for \( f \in L^p \). The right hand side of (2.12) is the Marchaud fractional derivative of \( \psi_x(u) \) (see (2.6)) evaluated for \( u = 1 \).

Proof. Let \( \alpha = k/2 \). As in [8, pp. 159, 161] we have

\[
\sum_{j=0}^{\ell} \binom{\ell}{j} (-1)^j (I_{-}^{\alpha}g_x)(1 + j t) = t^\alpha \int_0^\infty k(u)g_x(1 + ut) du,
\]

\[
k(u) = \frac{1}{\Gamma(\alpha)} \sum_{j=0}^{\ell} \binom{\ell}{j} (-1)^j (u - j)^{2-\alpha},
\]

and therefore

\[
(\mathbb{D}_x^\alpha \psi_x)(1) = (\mathbb{D}_x^\alpha I_{-}^{\alpha}g_x)(1) = \int_0^\infty \lambda_{t, \alpha}(\eta) g_x(1 + \varepsilon \eta) d\eta,
\]

\[
(\mathbb{D}_x^\alpha \psi_x)(1) = (\mathbb{D}_x^\alpha I_{-}^{\alpha}g_x)(1) = \int_0^\infty \lambda_{t, \alpha}(\eta) g_x(1 + \varepsilon \eta) d\eta.
\]
\[ \lambda_{\ell,a}(\eta) = [\kappa_\ell(\alpha) \Gamma(1 + \alpha)\eta]^{-1} \sum_{j=0}^{\ell} \binom{\ell}{j} (-1)^j (\eta - j)^\alpha. \]

It is known [8, Lemma 10.17] that
\[ \int_0^\infty \lambda_{\ell,a}(\eta) d\eta = 1, \quad \lambda_{\ell,a}(\eta) = \begin{cases} O(\eta^{\alpha-1}), & \eta < 1, \\ O(\eta^{\alpha-\ell-1}), & \eta > 1. \end{cases} \]

Owing to (2.7), (2.13) yields
\[ (\mathbb{D}^\alpha_{x}(\psi))(1) = \frac{2}{\sigma_{n-1}} \int_X f(y) k_\varepsilon([x,y]) dy, \quad k_\varepsilon(\tau) = \frac{(\tau^2 - 1)^{1-n/2}}{\varepsilon} \lambda_{\ell,a}\left(\frac{\tau^2 - 1}{\varepsilon}\right), \]
and application of Theorem 2.3 (on approximate identity) from [2] gives the required result. \(\square\)

Theorem A is a particular case of Theorem 2.2 for \(k = 1\).

2.2. The case \(X = S^n\). In this case an analog of (2.6) reads
\[ (I_{0+}^{k/2} g_x)(u) = \psi_x(u), \]

\[ g_x(v) = v^{-1/2}(M_{\sqrt{x}}(x)), \quad \psi_x(u) = \frac{u^{(k-1)/2}}{2\pi^{k/2}} \varphi \cos^{-1}\sqrt{x}(u), \]

with the left-sided (!) fractional integral \((I_{0+}^\alpha g)(u) = \frac{1}{\Gamma(\alpha)} \int_0^u g(v)(u - v)^{\alpha-1} dv\) and
\[ (M_{\sqrt{x}}(x)) = \frac{(1 - t^2)^{(1-n)/2}}{\sigma_{n-1}} \int_{\{y \in S^n : x \cdot y = t\}} f(y) d\sigma(y), \]
\(x \cdot y = x_1y_1 + \ldots + x_{n+1}y_{n+1}\). Now we proceed as in Section 2.1 with the following changes: we extend both sides of (2.14) to \(u < 0\) by zero, and then apply the “left-sided” Marchaud’s fractional derivative
\[ (\mathbb{D}^\alpha_{x} \psi)(u) = \lim_{\varepsilon \to 0} \frac{1}{\kappa_\ell(\alpha)} \int_{\varepsilon}^{\infty} \left[ \sum_{j=0}^{\ell} \binom{\ell}{j} (-1)^j \psi(u - j \varepsilon) \right] \frac{dt}{t^{1+\alpha}}. \]

By making use of approximation to the identity (see [11, Lemma 2.2] or [9, Lemma 2.9]) we get the following analog of Theorem 2.2.

**Theorem 2.4.** Let \(\varphi = \hat{f}, f \in L^p(S^n), 1 \leq p < \infty\). Then
\[ f(x) = \frac{\pi^{-k/2}}{2\kappa_\ell(k/2)} \int_0^\infty \left[ \sum_{j=0}^{\ell} \binom{\ell}{j} (-1)^j (1 - j \varepsilon)^{(k-1)/2} \varphi \cos^{-1}\sqrt{1-j \varepsilon}(x) \right] \frac{dt}{t^{1+k/2}}, \]
\[ \int_0^\infty = \lim_{\varepsilon \to 0} \int_{\varepsilon}^\infty, \]
where the limit is understood in the \(L^p\)-norm and in the a.e. sense. If \(f \in C(S^n)\), then \(\lim_{\varepsilon \to 0}\) can be treated in the sup-norm.

Theorem B is a particular case of Theorem 2.4 for \(k = 1\).
3. Inversion Formulae with Derivatives

We restrict ourselves by the case $X = \mathbb{H}^n$ (the case $X = S^n$, which was studied in [9], is simpler). Fractional integrals $\psi_x(u) = (I_{-}^{k/2}g_x)(u)$ in (2.6) are $L^q$-valued functions of $u$ for some $q$, and their derivatives

$$\frac{d}{du}\psi_x(u) = \lim_{\delta \to 0} \frac{\psi_x(u + \delta) - \psi_x(u)}{\delta}$$

can be defined in a natural way by setting $\lim = \lim_{\sigma \to 0}$ or $\lim = \lim$ (in the $x$-variable).

In accordance with this definition we have the following

**Lemma 3.1.** Let $g_x$ be defined by (2.7) with $f \in L^p$.

(i) If $1 \leq p < n - 1$, then $(-\frac{d}{du})(I_{-}^{\alpha}g_x)(u)|_{u=1} = f(x)$, the derivative being interpreted in the $L^p$-norm or in the a.e. sense.

(ii) If $0 \leq 1/q < 1/p - (2\alpha - 3)/(n - 1)$, then for each $u \geq 1$, $(-\frac{d}{du})(I_{-}^{\alpha}g_x)(u) = (I_{-}^{\alpha-1}g_x)(u)$ in the $L^q$-norm or in the a.e. sense.

**Proof.** (i) For $\psi = I_{-}^{\alpha}g_x$ the required derivative can be expressed through approximate identity. Namely,

$$\lim_{\delta \to 0} \frac{\psi(1) - \psi(1 + \delta)}{\delta} = \lim_{\delta \to 0} \frac{1}{\delta} \int_{1}^{1+\delta} M_x f \frac{dx}{\sqrt{\tau}} = \frac{2}{\sigma_{n-1}} \lim_{\delta \to 0} \int_X k_\delta([x,y])f(y)dy,$$

$$k_\delta(\tau) = \frac{(\tau^2 - 1)^{1-n/2}}{\delta} \chi \left( \frac{\tau^2 - 1}{\delta} \right), \quad \chi(s) = \begin{cases} 1 & \text{if } 0 < s < 1, \\ 0 & \text{if } s > 1, \end{cases}$$

and the result follows by Theorem 2.3 from [2].

(ii) We have to show that the expression $\frac{1}{\delta} \left[ (I_{-}^{\alpha}g_x)(u) - (I_{-}^{\alpha}g_x)(u+\delta) - (I_{-}^{\alpha-1}g_x)(u) \right]$ tends to zero as $\delta \to 0$ in the required sense. It can be written as a convolution

$$\frac{2}{\sigma_{n-1}}(f * h_\delta)(x)$$

where

$$h_\delta(\tau) = (\tau^2 - 1)^{1-n/2} \left[ \frac{(\tau^2 - u)^{\alpha-1}}{\Gamma(\alpha)} - \frac{(\tau^2 - u - \delta)^{\alpha-1}}{\Gamma(\alpha)} \right]$$

$$= \frac{(\tau^2 - 1)^{1-n/2}(\tau^2 - u)^{\alpha-2}}{\Gamma(\alpha-1)} \cdot \frac{\delta}{\tau^2 - u} - k(\omega) = \frac{1 - (1 - \omega)^{\alpha-1}}{(\alpha - 1)\omega} - 1.$$

The function $k(\omega)$ is bounded for $\omega > 0$ and $\lim_{\omega \to 0} k(\omega) = 0$. Hence by Young’s inequality, $\|f * h_\delta\|_q \leq \|f\|_p\|h_\delta\|_r$ with the required $q$ (use (2.9) with $\alpha$ replaced by $\alpha - 1$). Now the Lebesgue theorem on dominated convergence yields

$$\lim_{\delta \to 0} \|h_\delta\|_r^a = 0, \quad \lim_{\delta \to 0} (f * h_\delta)(x) = 0,$$
and we are done.

Lemma 3.1 leads to the following

**Definition 3.2.** We set

\[
D_m = (-1)^m \partial_m \partial_{m-1} \ldots \partial_1,
\]

where all derivatives \( \partial_i \) designate the same expression \( \frac{d}{du} \), but their interpretation depends on \( i \). Namely, if \( i < k/2 \), then \( \partial_i \) is understood in the \( L^\infty \)-norm where

\[
0 \leq \frac{1}{q_i} < \frac{1}{p} - \frac{k - 1}{n - 1} + \frac{2i}{n - 1}.
\]

If \( i = k/2 \), then \( \partial_i \) is understood in the \( L^p \)-norm. Alternatively all derivatives in (3.1) may be treated in the a.e. sense.

**Theorem 3.3.** Let \( \varphi = \tilde{f} \), \( f \in L^p \), \((k - 1)/(n - 1) < 1/p \leq 1\),

\[
\tilde{\psi}_m(x,u) = \pi^{-k/2} D_m [u^{(k-1)/2} \varphi \cosh^{-1} \sqrt{u} (x)], \quad m \in \mathbb{N}.
\]

(i) If \( k = 2m \), then \( f(x) = \tilde{\psi}_m(x,1) \).

(ii) If \( k = 2m + 1 \), then

\[
f(x) = \frac{1}{2\pi^{1/2}} \int_0^\infty \left[ \tilde{\psi}_m(x,1) - \tilde{\psi}_m(x,1 + t) \right] \frac{dt}{t^{3/2}}, \quad \int_0^\infty = \lim_{\varepsilon \to 0} \int_0^\infty ,
\]

where the limit is understood as in Theorem 2.2.

**Proof.** (i) follows immediately from Lemma 3.1 and Definition 3.2. In order to prove (ii), by Lemma 3.1 we have \((L^{1/2} g_x)(u) = \tilde{\psi}_m(x,u)\). As in the proof of Theorem 2.2, application of Marchaud’s derivative \( \mathbb{D}^{1/2} \) yields

\[
f(x) = \frac{1}{\varpi_1(1/2)} \int_0^\infty \left[ \tilde{\psi}_m(x,1) - \tilde{\psi}_m(x,1 + t) \right] \frac{dt}{t^{3/2}}, \quad \varpi_1(1/2) = 2\pi^{1/2}.
\]

This completes our investigation. \( \Box \)

**References**


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