

*Happy Birthday  
Yosi !*

# MODELING LIQUID METALS and BULK METALLIC GLASSES

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# Content

1. Metal Liquids and Glasses
2. Delone Graphs
3. Conclusion

# I - Metal Liquids and Glasses

# Bulk Metallic Glasses

## 1. Examples *(Ma, Stoica, Wang, Nat. Mat. '08)*

- $\text{Zr}_x\text{Cu}_{1-x}$      $\text{Zr}_x\text{Fe}_{1-x}$      $\text{Zr}_x\text{Ni}_{1-x}$
- $\text{Cu}_{46}\text{Zr}_{47-x}\text{Al}_7\text{Y}_x$      $\text{Mg}_{60}\text{Cu}_{30}\text{Y}_{10}$

## 2. Properties *(Hufnagel web page, John Hopkins)*

- High *Glass Forming Ability* (GFA)
- High *Strength*, comparable or larger than steel
- Superior *Elastic limit*
- High *Wear* and *Corrosion* resistance
- *Brittleness* and *Fatigue* failure

# Bulk Metallic Glasses

Applications (Liquidmetal Technology [www.liquidmetal.com](http://www.liquidmetal.com))

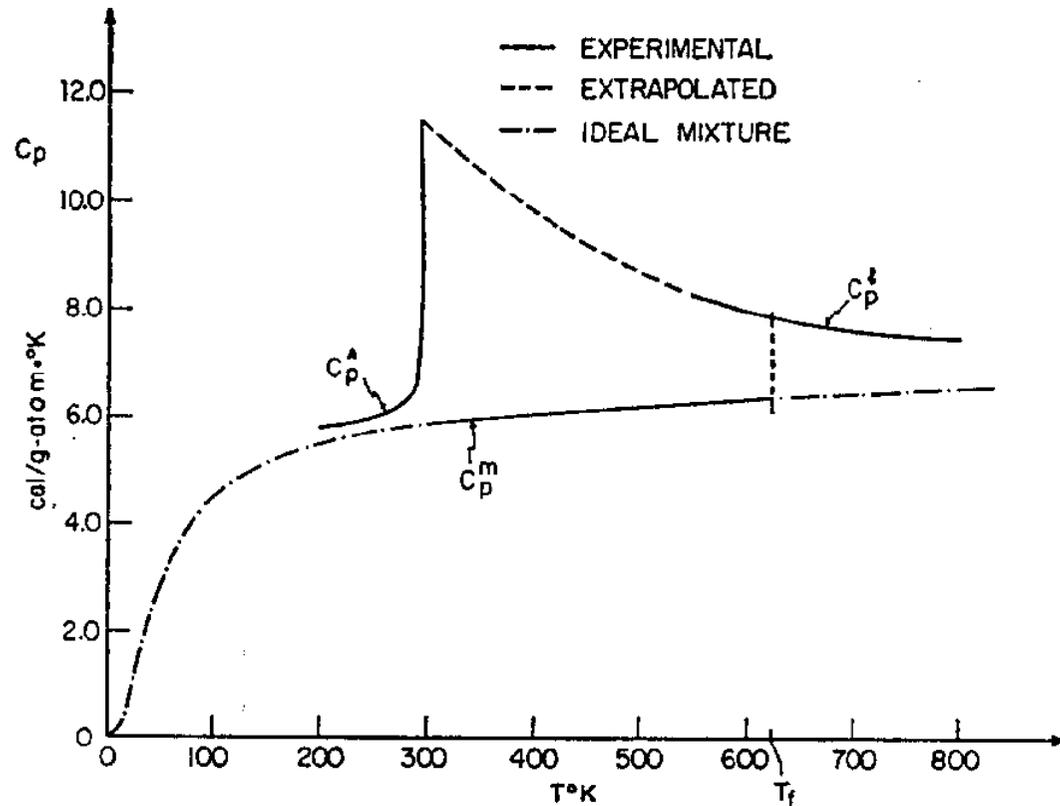
- *Orthopedic implants* and medical Instruments
- Material for *military components*
- Sport items, *golf clubs, tennis rackets, ski, snowboard, ...*



Pieces of Titanium-Based Structural  
Metallic-Glass Composites

(Johnson's group, Caltech, 2008)

# Bulk Metallic Glasses

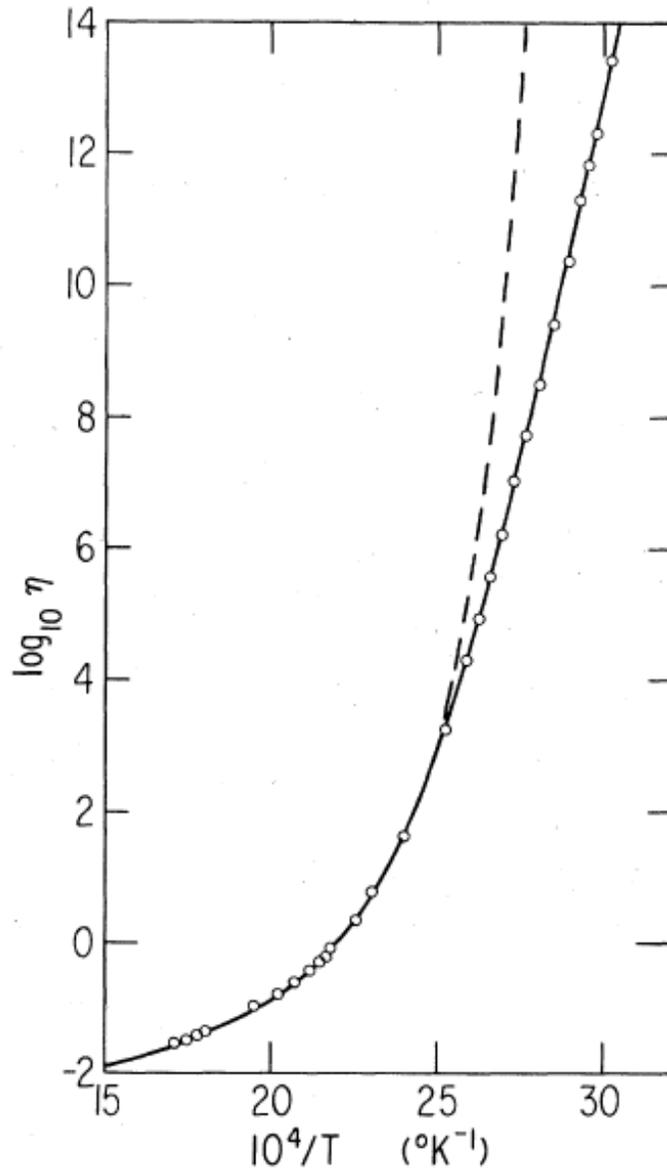


Smoothed values of specific heats of  $Au_{.77}Ge_{.136}Si_{.094}$  signaling a glass-liquid transition

"A" designates the amorphous state  
"m" designates the mixture  
"l" designates the liquid

taken from  
H. S. CHEN and D. TURNBULL, *J. Chem. Phys.*,  
48, 2560-2571, (1968)

# Bulk Metallic Glasses



Viscosity vs temperature for tri-anaphthylbenzene, with fits coming from the *free volume theory*

*Solid curve* fit from [1] below

*Dashed curve*: fit from [1] with a simplified theory

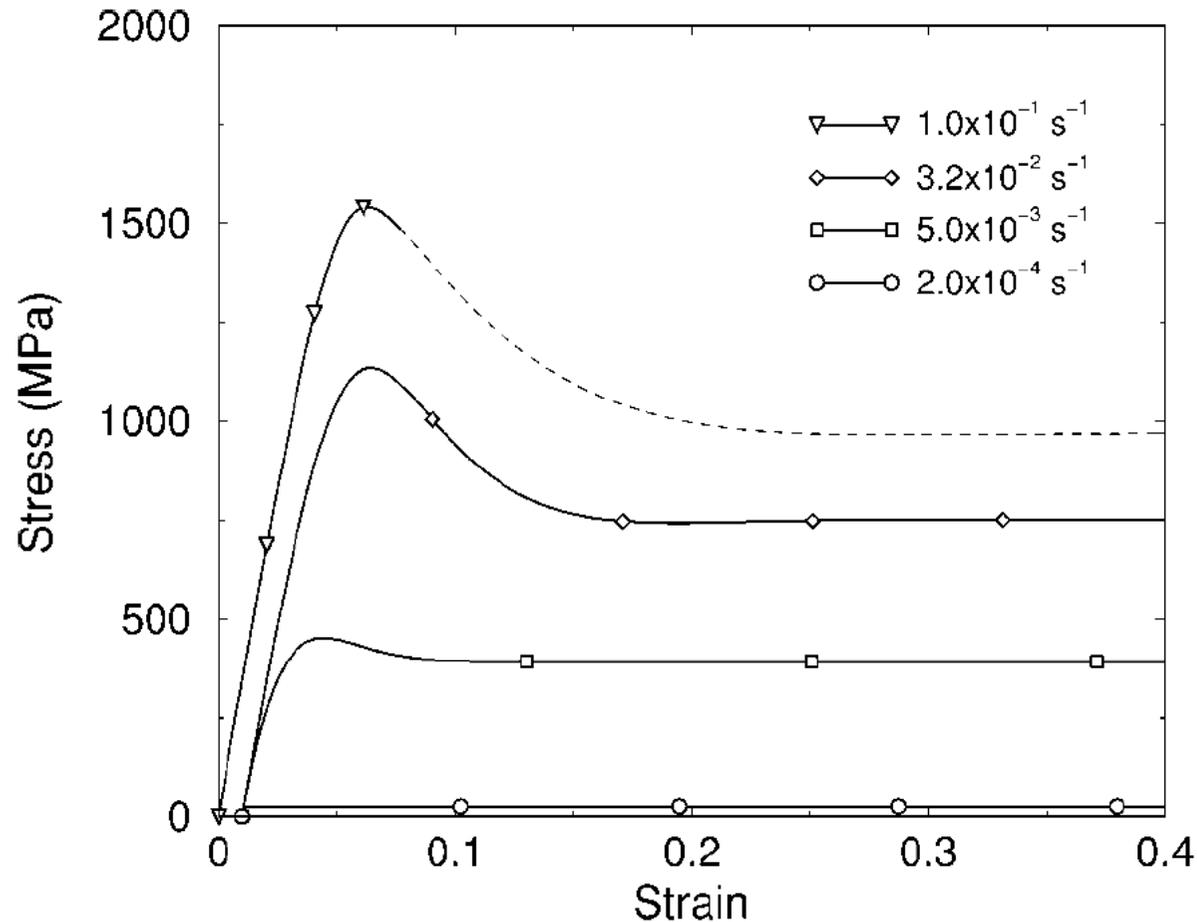
*Circles*: data from [2] below

*taken from*

[1] MORREL H. COHEN & G. S. GREST, *Phys. Rev. B*, **20**, 1077-1098, (1979)

[2] D. J. PLAZEK and J. H. MAGILL, *J. Chem. Phys.*, **45**, 3757, (1967); J. H. MAGILL, *ibid.* **47**, 2802, (1967)

# Bulk Metallic Glasses



Theoretical curves of tensile stress versus strain for the bulk metallic glass using the *STZ theory*



at several different strain rates as shown. The temperature is  $T=643 \text{ K}$ .

*For clarity, all but the first of these curves have been displaced by the same amount along the strain axis.*

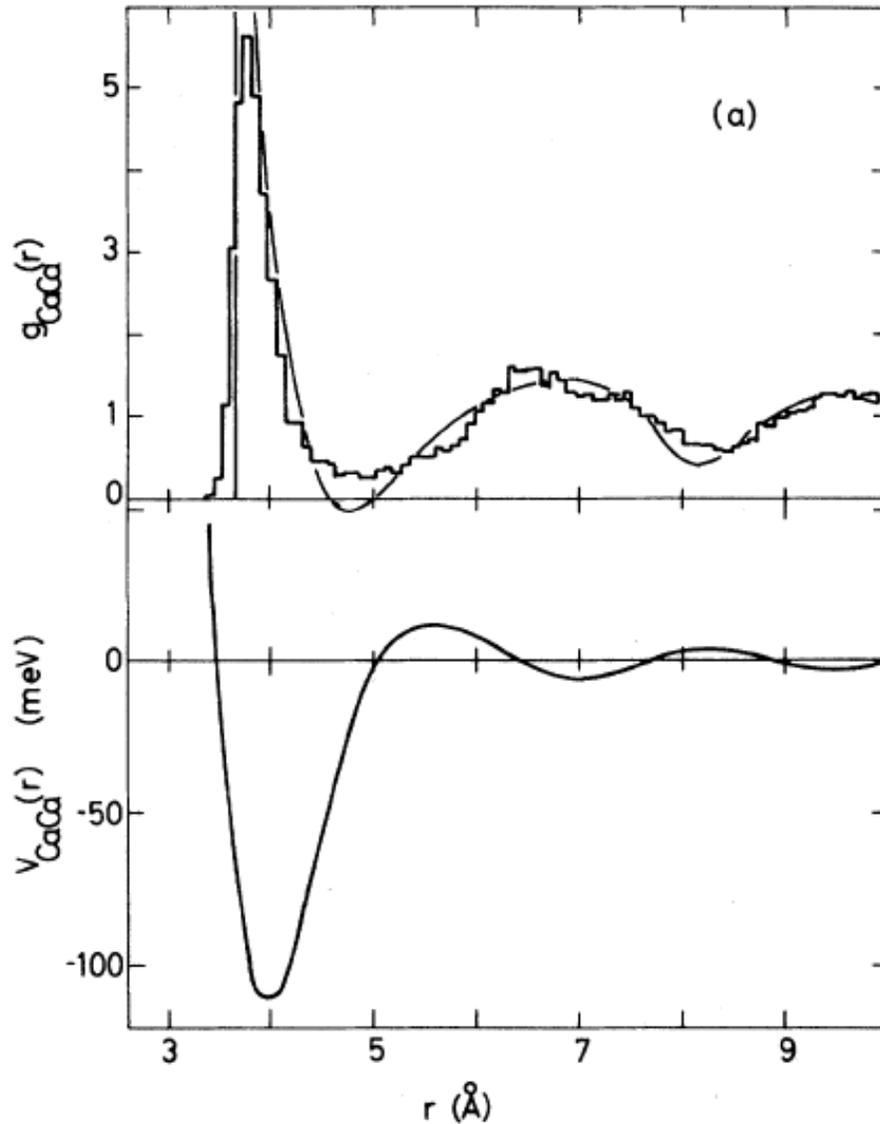
*taken from*

[1] M. L. FALK, J. S. LANGER & L. PECHENIK, *Phys. Rev. E*, **70**, 011507, (2004)

# Available Theories

1. **Cluster Model:** (*T. Egami et al., '80's, D. Miracle '04*) gives a guide line to how to produce glassy state, explains the diffraction spectrum and the pair-distribution function. It introduces *frustration*
2. **Free Volume Theory:** (*Morell H. Cohen, G.S. Grest '79*) gives a good account for the viscosity
3. **Shear Transformation Zone (STZ):** (*J. Langer et al. '98-06*) gives effective continuum equation valid beyond the elasticity limit. Numerical simulations give a convincing description of the dynamics of cracks (*C.H. Rycroft, E. Bouchbinder '12*)

# Pair Potentials



An example of atom-atom pair potential in the metallic glass  $\text{Ca}_{70}\text{Mg}_{30}$

*Top:* the pair creation function  
*Bottom:* the graph of the pair potential

*taken from*  
J. HAFNER, *Phys. Rev. B*, 27, 678-695 (1983)

# Dense Packing

1. The shape of the pair potential suggests that there is a *minimal distance* between two atoms.
2. Liquid and solids are *densely packed*. This suggests that there is a *maximal size for vacancies*.
3. However, Mathematics (*ergodic theory*) implies that, given an  $\epsilon > 0$ , with probability one
  - there are pairs of atoms with distance less than  $\epsilon$
  - there are vacancies with radius larger than  $1/\epsilon$
4. But these rare events are not seen in practice because their *lifetime is negligibly small* (Bennett et al. '79). Persistence theory give an argument in this direction.

## II - Delone Graphs

# Delone Sets

- The set  $\mathcal{V}$  of atomic positions is *uniformly discrete* if there is  $b > 0$  such that in any ball of radius  $b$  there is at *most* one atomic nucleus.

(Then minimum distance between atoms is  $\geq 2b$ )

- The set  $\mathcal{V}$  is *relatively dense* if there is  $h > 0$  such that in any ball of radius  $h$  there is at *least* one atomic nucleus.

(Then maximal vacancy diameter is  $\leq 2h$ )

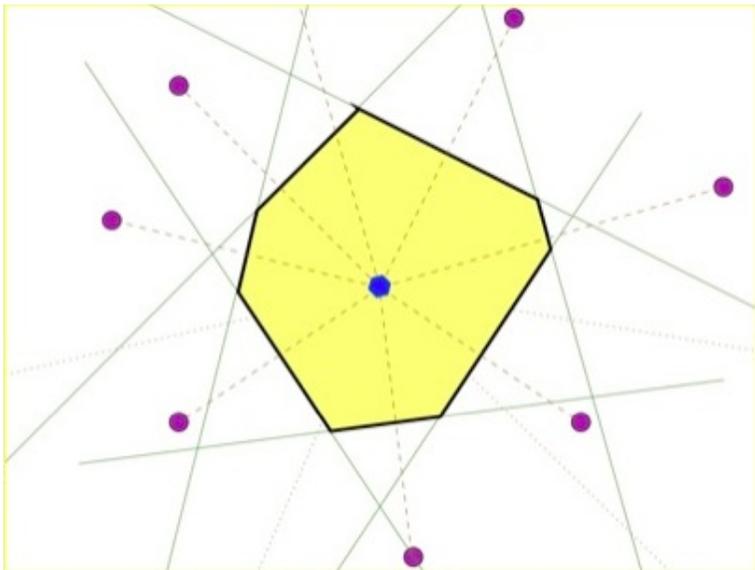
- If  $\mathcal{V}$  is both uniformly discrete and relatively dense, it is called a *Delone set*.
- $\text{Del}_{b,h}$  denotes the set of *Delone sets* with parameters  $b, h$ .

# Voronoi Cells

- Let  $\mathcal{V} \in \text{Del}_{b,h}$ . If  $x \in \mathcal{V}$  its *Voronoi cell* is defined by

$$V(x) = \{y \in \mathbb{R}^d ; |y - x| < |y - x'| \forall x' \in \mathcal{V}, x' \neq x\}$$

$V(x)$  is open. Its closure  $T(x) = \overline{V(x)}$  is the *Voronoi tile* of  $x$



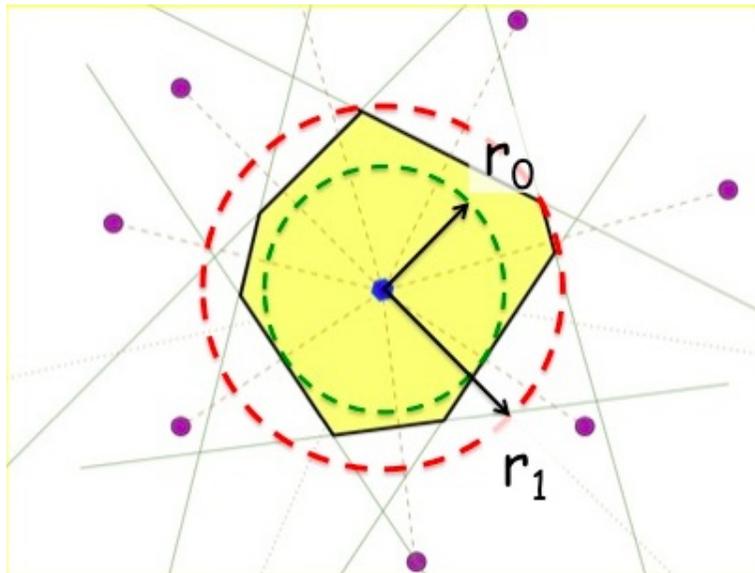
**Proposition:** If  $\mathcal{V} \in \text{Del}_{r_0,r_1}$  the Voronoi tile of any  $x \in \mathcal{V}$  is a convex polytope

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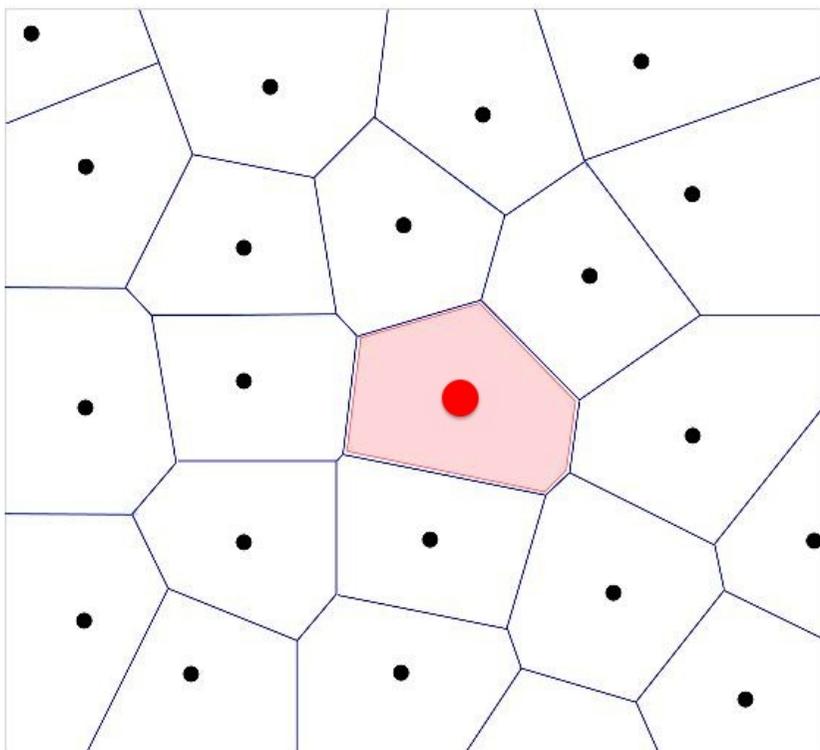
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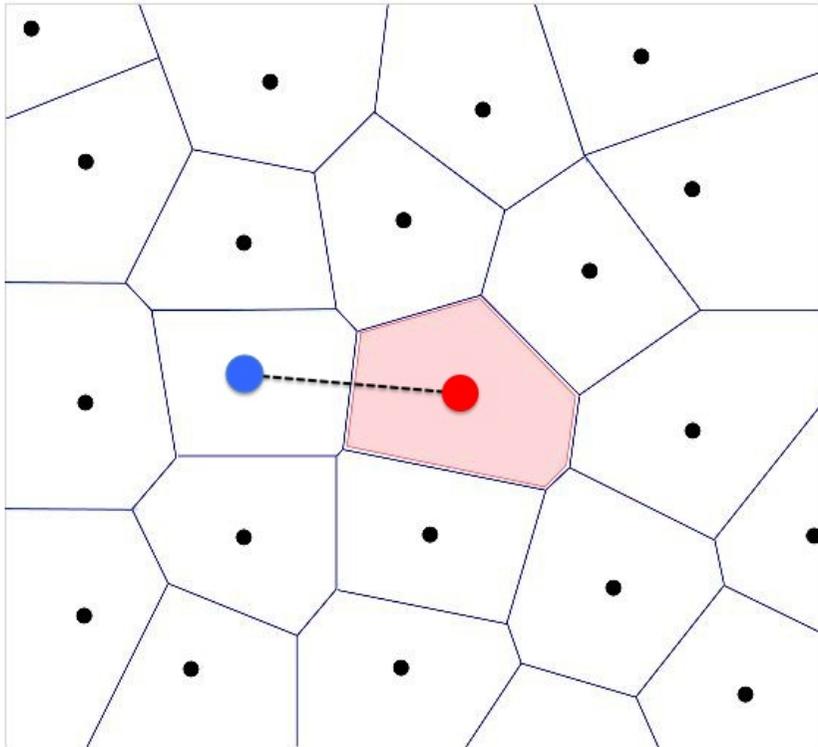
**Proposition:** If  $\mathcal{V} \in \text{Del}_{r_0, r_1}$  the Voronoi tile of any  $x \in \mathcal{V}$  is a convex polytope containing the ball  $\overline{B}(x; r_0)$  and contained in the ball  $\overline{B}(x; r_1)$

# The Delone Graph



**Proposition:** *the Voronoi tiles of a Delone set touch face-to-face*

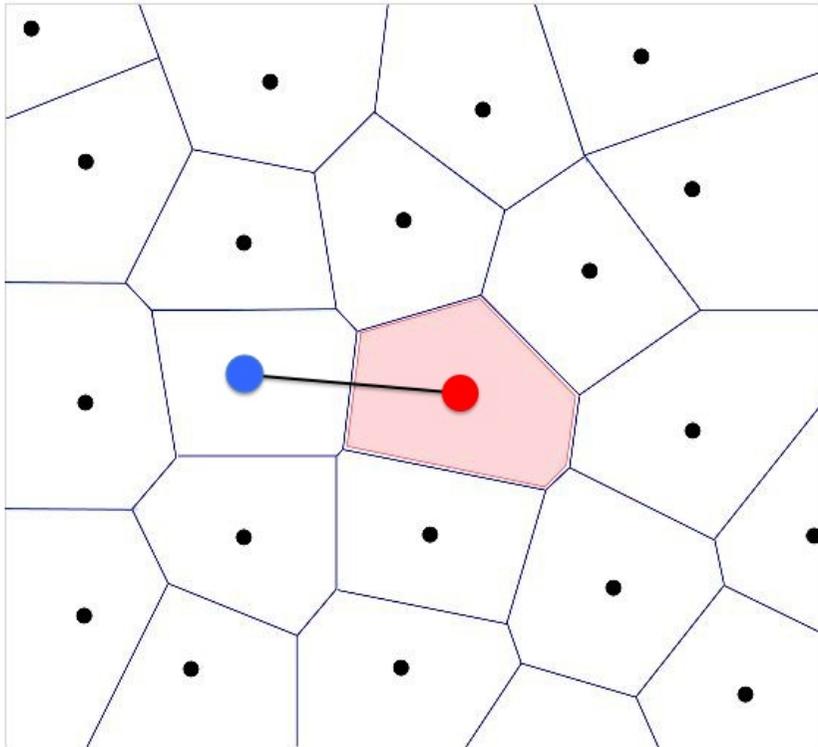
# The Delone Graph



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Two atoms are *nearest neighbors* if their Voronoi tiles touch along a face of *maximal dimension*.

# The Delone Graph

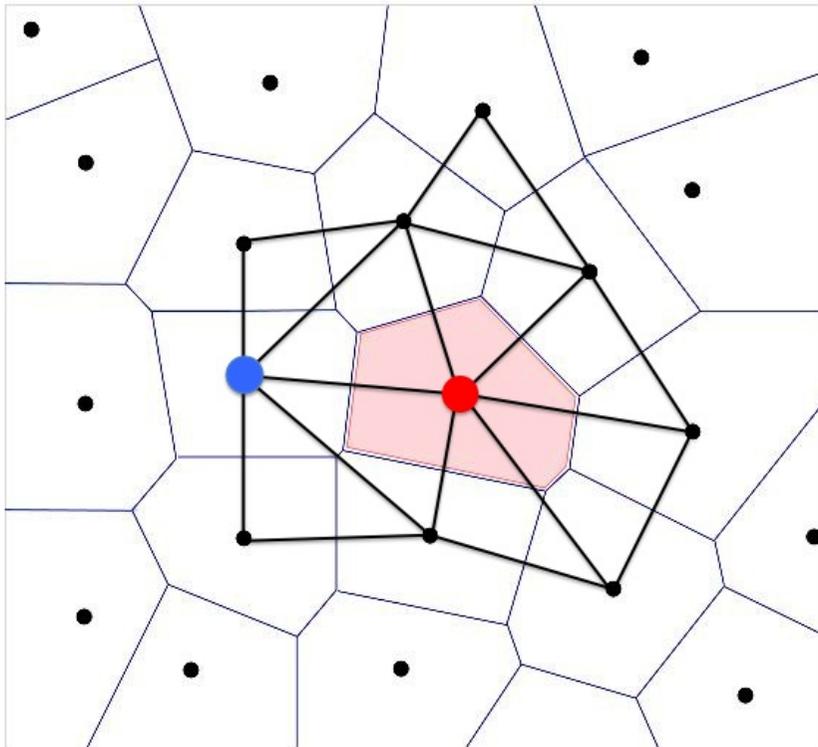


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An *edge* is a pair of nearest neighbors.  $\mathcal{E}$  denotes the set of edges.

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The family  $\mathcal{G} = (\mathcal{V}, \mathcal{E})$  is the Delone graph.

# The Delone Graph

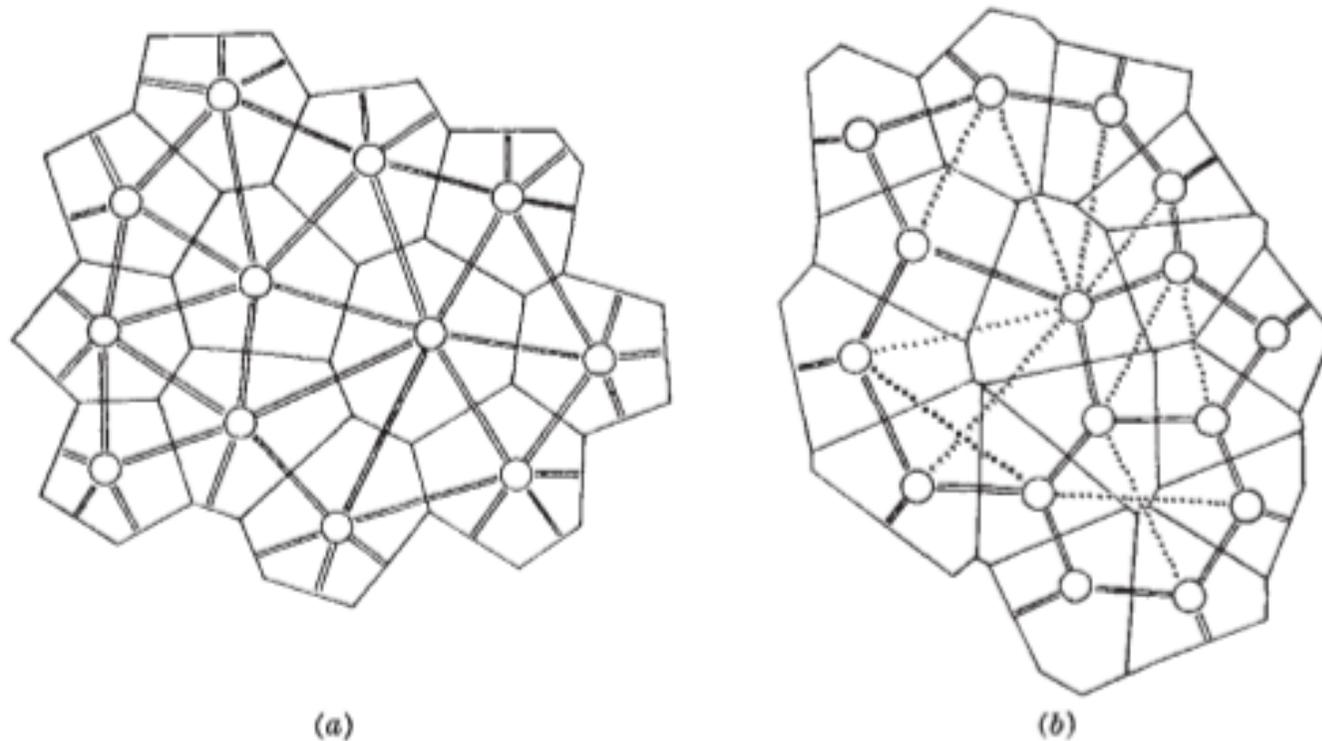


Fig. 1. Diagram of neighbourhood polyhedra, geometrical and physical, for two-dimensional arrays of points. (a) High co-ordinated; =, physical neighbours; (b) low co-ordinated; ....., geometrical neighbours

*taken from J. D. BERNAL, Nature, 183, 141-147, (1959)*

# Properties of the D-graph

- A Delone graph is *simple*: one edge at most between two vertices, no edge with one end point (tadpole).
- A *graph map* sends vertices to vertices, edge to edges and is compatible with the edge boundaries
- Graphs are identified modulo *graph isomorphisms*
- *Given an integer  $N$ , the number of simple graphs modulo isomorphism with less than  $N$  vertices is finite*
- **Consequence:** There are only finitely many D-Graphs representing a configuration of the glass in a ball of finite radius. D-graphs *discretize* the information.

# Properties of the D-graph

- The incidence number  $n_x$  of a vertex  $x \in \mathcal{V}$  is bounded by

$$d + 1 \leq n_x \leq \frac{\sqrt{\pi} \Gamma\{(d - 1)/2\}}{\Gamma(d/2) \int_0^{\theta_m} \sin^{d-1}(\theta) d\theta}, \quad \sin \theta_m = b/2h.$$

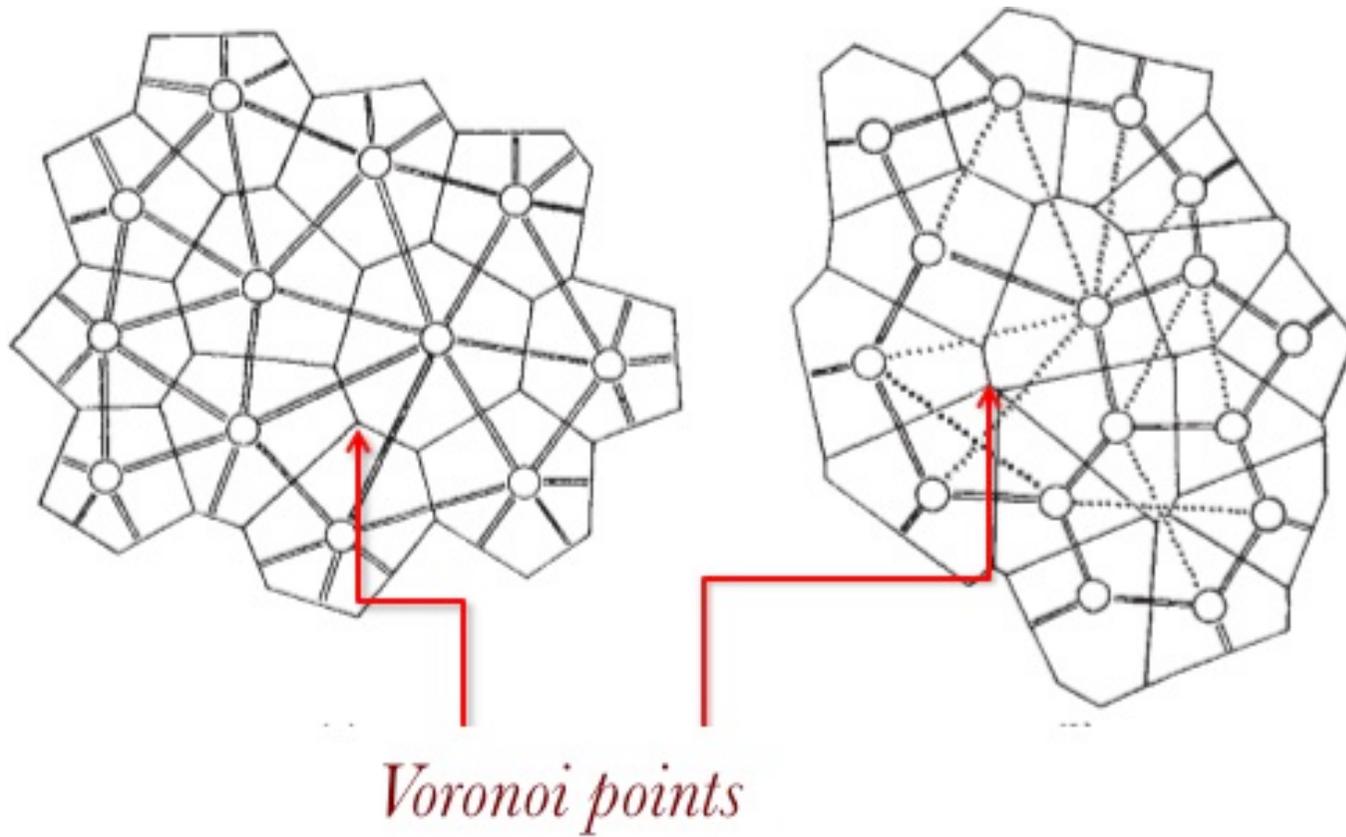
- A *local patch* of radius  $n \in \mathbb{N}$  is an *isomorphism class* of subgraphs  $(x, \mathcal{V}_x, \mathcal{E}_x)$  of the Delone graph, such that  $x \in \mathcal{V}$ ,  $\mathcal{V}_x$  is the set of vertices at graph-distance at most  $n$  from  $x$ .
- If  $\mathcal{P}_n$  denote the *set of local patches* of radius  $n$  then there is  $C = C(b, h) > 0$  such that

$$\#\mathcal{P}_n \leq e^{C(2n+1)^d}$$

# Likelihood

- Likelihood can be expressed in various ways such as *topological genericity* or *full measure* (say *w.r.t* the Lebesgue measure)
- *If  $X \subset \mathbb{R}^n$  is closed and if  $\mathbb{P} = F(x) d^n x$  is “absolutely continuous”, then a property valid of a dense open set  $U \subset X$ , with piecewise smooth boundary, is both generic and almost sure.*

# Voronoi Points

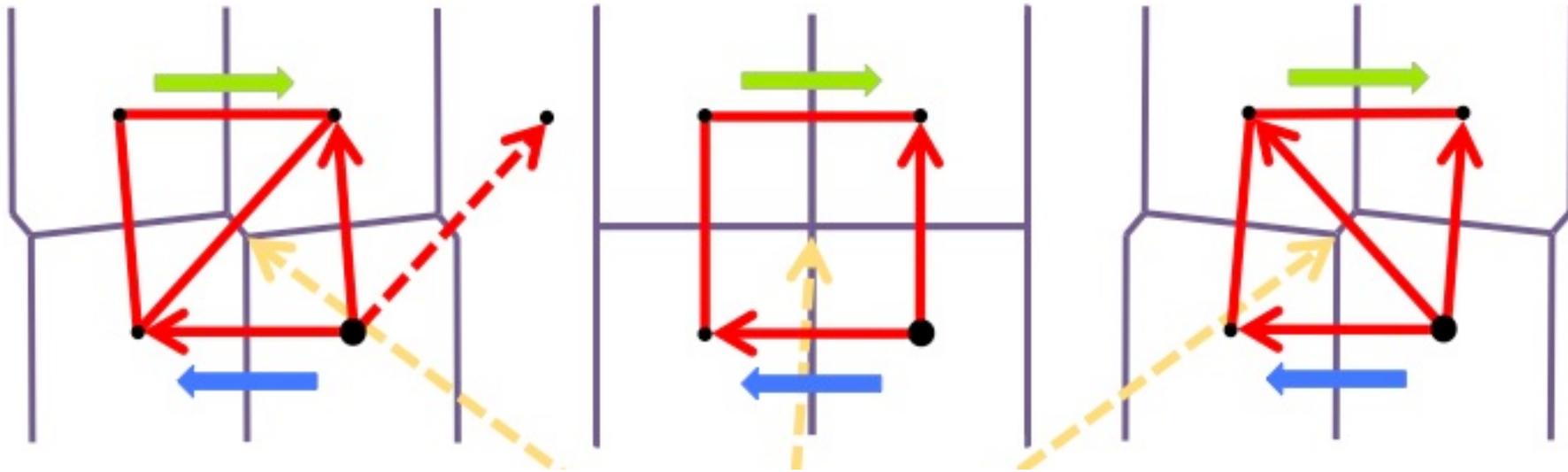


The vertices of the Voronoi cells are called *Voronoi Points*.

# Voronoi Points

- **Theorem:** *A Voronoi point is at equal distance from every atom the Voronoi tile of which it belongs to.*
- **Theorem:** *Generically and almost surely a Voronoi point belongs to exactly  $d + 1$  Voronoi tiles in dimension  $d$ .*

# Generic Local Patches

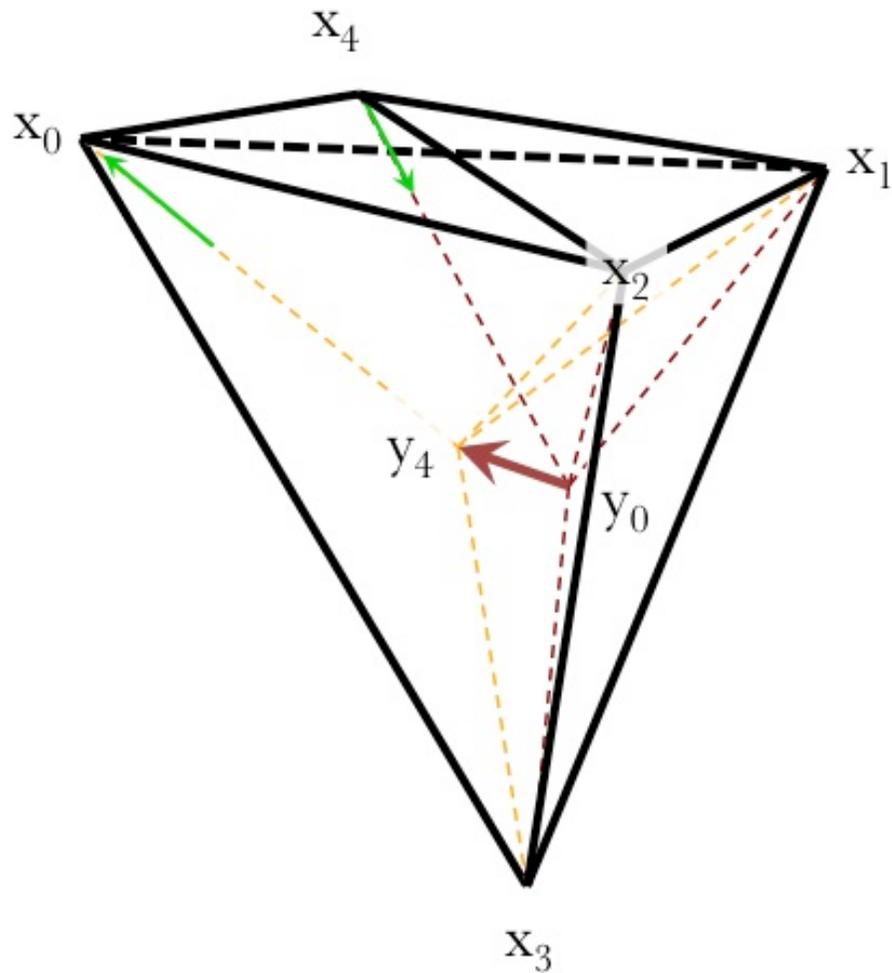


Shear modifies local patches. The middle one is *unstable*. The transition from left to right requires transiting through a *saddle point* of the potential energy.

*The Voronoi cell boundaries are shown in blue.*

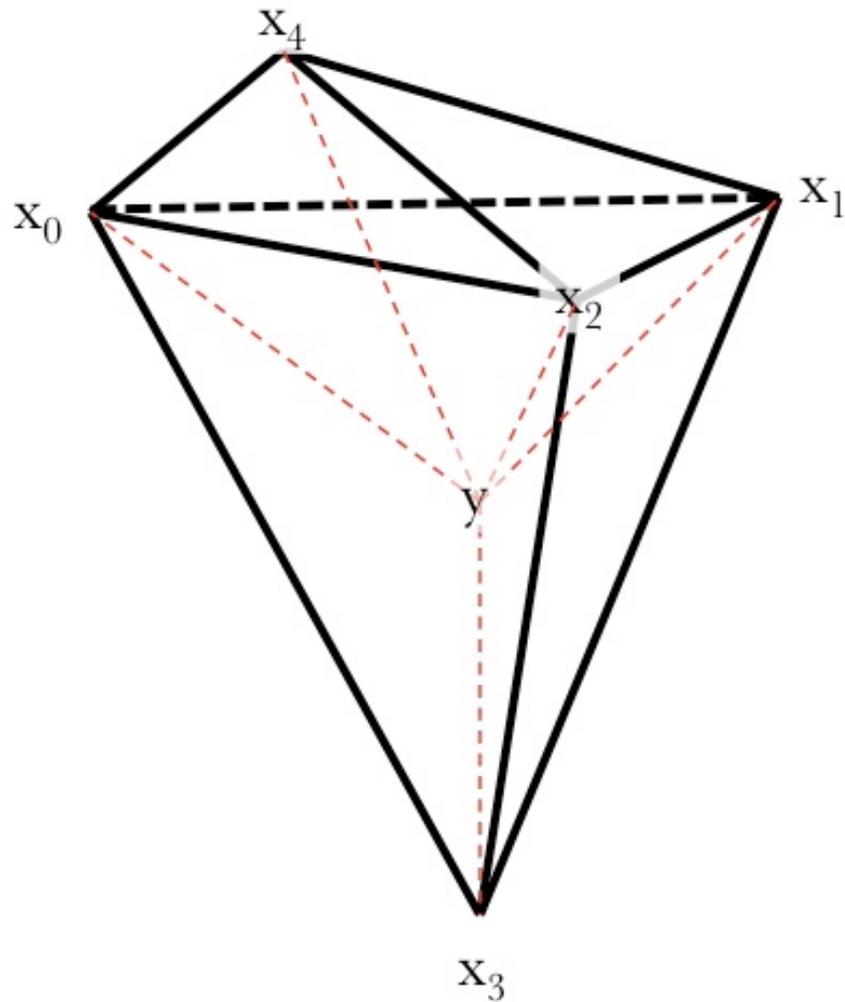
*At the bifurcation a Voronoi vertex touches one more Voronoi cell than in the generic case*

# Generic Local Patches



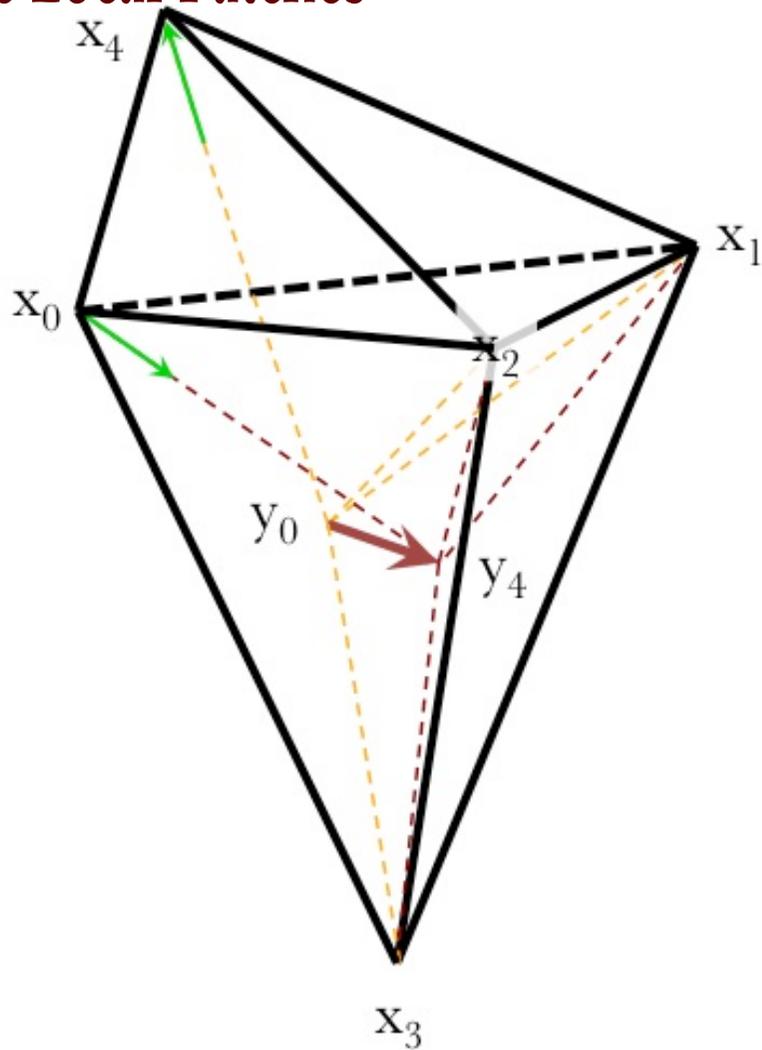
*An example of a generic 3D bifurcation.*

# Generic Local Patches



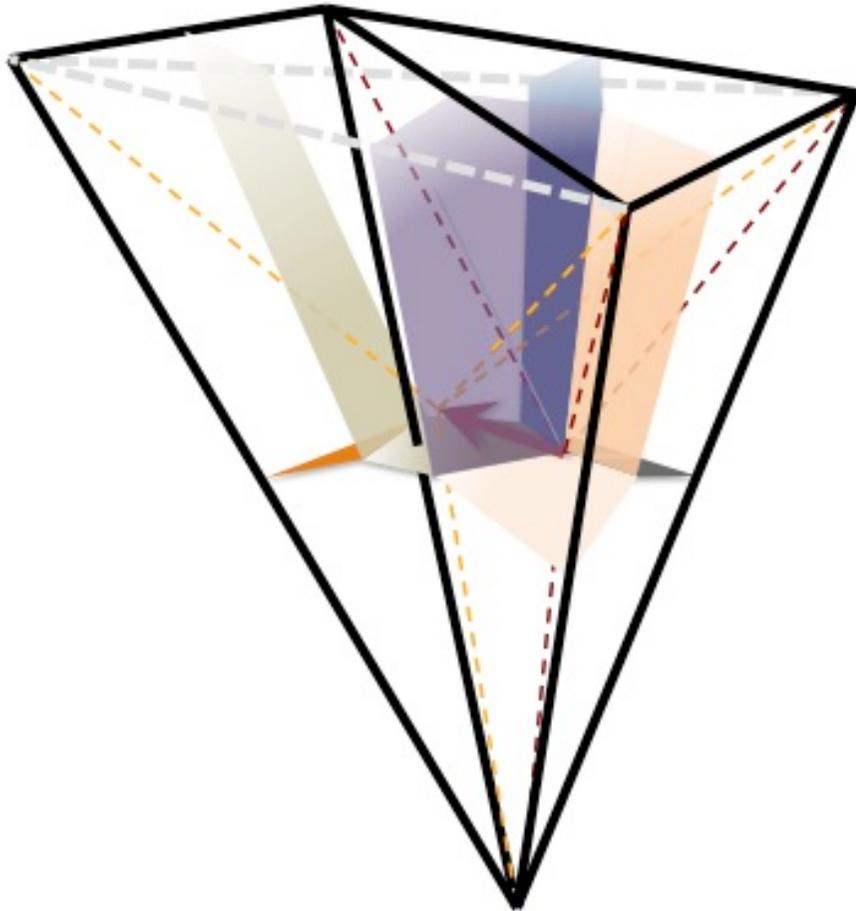
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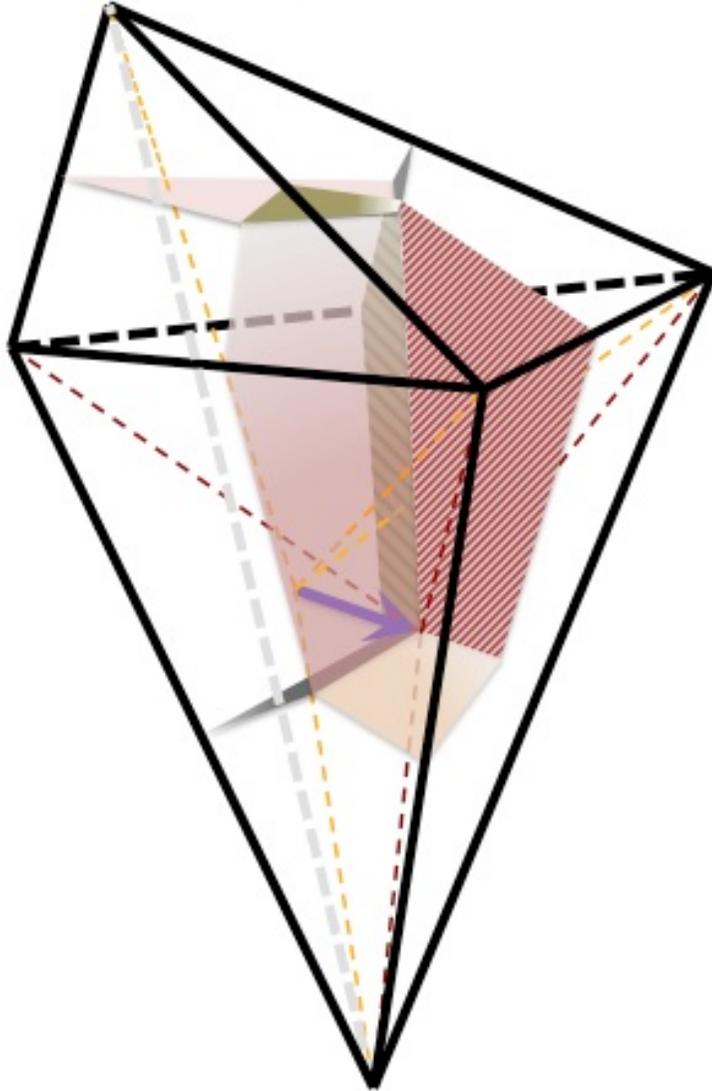


*An example of a generic 3D bifurcation.*

*Graph changes*

- The graph edges are indicated in black.
- The grey dotted edges have disappeared during the bifurcation.
- The colored plates are the boundaries of the Voronoi cells.

# Generic Local Patches



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*Graph changes*

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# Acceptance Domains

- Given a local patch  $\mathcal{G} \in \mathcal{P}_n$  its acceptance domain  $\Sigma_{\mathcal{G}}$  is the set of all atomic configurations  $\mathcal{V} \in \text{Del}_{b,h}$  having  $\mathcal{G}$  as their *local patch around the origin*.
- A local patch is *generic* whenever a small local deformation of the atomic configuration does not change the corresponding graph. Let  $\mathfrak{B}_n \subset \mathcal{P}_n$  be the set of *generic local patches* of radius  $n$ .
- **Theorem:**  $\mathcal{G} \in \mathfrak{B}_n$  if and only if  $\Sigma_{\mathcal{G}}$  is open and its boundary is *piecewise smooth*.  
*The union of acceptance domains of the generic patches of size  $n$  is dense.*  
*In particular almost surely and generically an atomic configuration admits a generic local patch.*

# Contiguosness

- The *boundary* of the acceptance domain of a generic graph contains a relatively open dense subset of codimension 1.
- **Definition:** *two generic graphs  $\mathcal{G}, \mathcal{G}' \in \mathcal{Q}_n$  are contiguous whenever their boundary share a piece of codimension one.*
- The set  $\mathcal{B}_n$  itself can then be seen as the set of vertices of a graph

$$\mathfrak{G}_n = (\mathcal{B}_n, \mathfrak{E}_n)$$

called the *graph of contiguousness* where *an edge  $E \in \mathfrak{E}_n$  is a pair of contiguous generic local patches.*

# Contiguousness

**Theorem** *two contiguous generic graphs differ only by one edge*

## III - Conclusion

# Degrees of Freedom

1. In the liquid phase, the short wave phonons are *short lived*. They do not contribute to the heat capacity.
2. The only relevant degree of freedom is the *bond motion*. The contiguousness relation expresses the combinatoric part of this degree of freedom
3. Each bond comes with a local *stress tensor*. In the liquid phase this stress tensor is *Gaussian distributed* (free Maxwell gas)  
⇒ *law of Dulong and Petit for the heat capacity*.
4. The bond degrees of freedom are called *anankeon*.

# The Anankeon Theory

*The bond degrees of freedom are the response of atoms to the **stressful** situation in which they are trying to find a better comfortable position, in vain.*

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**the goddess Ananke**

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*whose name comes from the greek word **anagkeia** meaning the **stress of circumstances**. Ananke was representing a power above all including the Gods of the Olympe "even gods don't fight against Ananke" claims a scholar. This character presided to the **creation of the world**, in various versions of the Greek mythology. It expresses the concepts of "**force, constraint, necessity**" and from there it also means "**fate, destiny**" to lead to the concepts of compulsion, torture.*(from Wikipedia)**

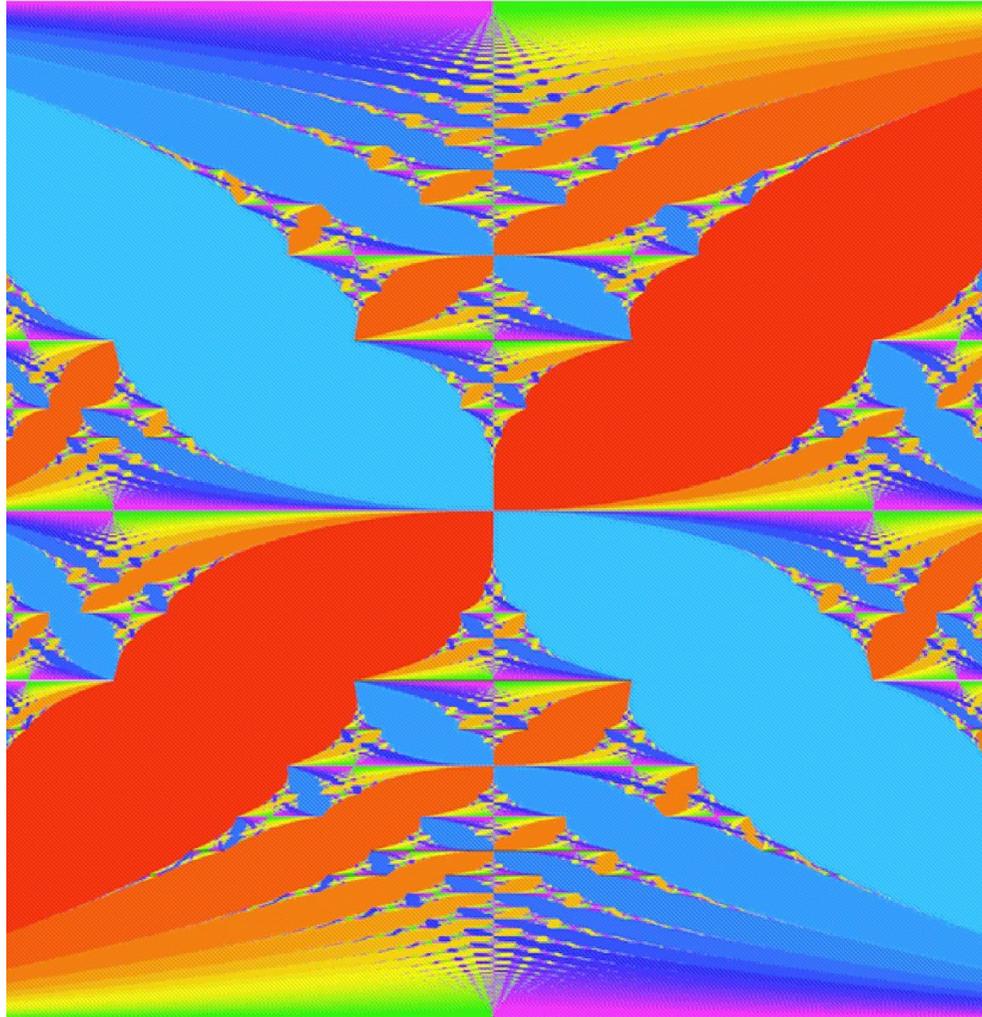
# The Anankeon Theory

For this reason the configurational degrees of freedom associated with the stress tensor on each bond will be called

**ANANKEONS**

# Modeling the Dynamics

1. A *Markov dynamics* can be built on the graph of *contiguosness*, describing the time evolution of the D-graph of both liquid and glasses, or the dynamics of the local bonds (*anankeons*)
2. Each D-graph is decorated by local degrees of freedom: local stress (*anankeon*), local vibration (*phonons*).
3. In the *liquid phase* the anankeon are free, phonons are suppressed, the theory gives a perfect gas solvable model.
4. The interaction *phonon-anankeon* is still ununderstood.



*Happy Birthday  
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